

# Some applications on tangent bundle with Kaluza-Klein metric

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**Abstract:** In this paper, differential equations of geodesics; parallelism, incompressibility and closeness conditions of the horizontal and complete lift of the vector fields are investigated with respect to Kaluza-Klein metric on tangent bundle.

**Keywords:** Kaluza-Klein metric, tangent bundle, geodesics, harmonic vector fields

## 1 Introduction

The history of tangent bundle  $TM$  of a Riemannian manifold  $M$  goes back with the work of Sasaki [1]. After the Dombrowski's comparison with the geometric objects of the base manifold and of the tangent bundle with respect to Sasaki metric  $^Sg$  in [2], many studies have been done on this topic in this sense and it has been seen that the metric  $^Sg$  arises a kind of "rigidity", i.e., most of geometric properties of the  $TM$  can not be ensured unless the base manifold  $M$  (and hence  $(TM, ^Sg)$ ) is flat. To eliminate this deficiency, Musso and Tricceri defined another metric and called as Cheeger-Gromoll metric  $^{CG}g$  [3]. It is shown that  $(TM, ^{CG}g)$  is not flat even if  $M$  is flat [4]. Later, more general metrics are demonstrated with deformations  $^{CG}g$  including both Sasaki and Cheeger-Gromoll metrics. One of them is the metric  $g_{a,b}$ , which is obtained from deformation of the vertical part of the metric  $^{CG}g$  with two positive definite scalar functions  $a$  and  $b$  [5].

In this paper, the Kaluza-Klein metric  $^{KK}g$ , which is defined by rescaling the horizontal part of the metric  $g_{a,b}$ , is considered and various characterizations on geodesics and some special vector fields are given with respect to this metric on tangent bundle. Throughout the paper; manifolds, functions and vector fields are differentiable of class  $C^\infty$ .

## 2 Preliminaries

Let  $M$  be a manifold with finite dimension  $n$ . Then the set  $TM = \bigcup_{P \in M} T_P M$  is the tangent bundle on  $M$ , where  $\cup$  denotes the disjoint union of the tangent spaces  $T_1^1(P)$  for all  $P \in M$ .  $TM$  is a  $2n$ -dimensional manifold. The natural projection  $\pi : T_1^1(M) \rightarrow M$  is defined for any point  $\tilde{P}$  of  $T_1^1(M)$  such that  $\tilde{P} \in T_1^1(P)$  with the surjective correspondence  $\tilde{P} \rightarrow P$ . If  $x^j$  are local coordinates in a neighbourhood  $U$  of  $P \in M$ , then a vector field  $u$  at  $P$  which is an element of  $TM$  is expressed in the form  $(x^j, u^j)$ , where  $u^j$  are components of  $t$  with respect to the natural base. We consider  $(x^j, u^j) = (x^j, x^{\bar{j}}) = (x^J)$ ,  $j = 1, \dots, n$ ,  $\bar{j} = n + 1, \dots, 2n$ ,  $J = 1, \dots, 2n$  as local coordinates in a neighborhood  $\pi^{-1}(U)$ .

Now, let we assume that the base manifold  $M$  is a Riemannian manifold with the metric  $g$  and denote by  $\nabla$  its Levi-Civita connection. Then the vertical distribution is  $VTM = \ker \pi_*$  and the horizontal distribution is defined by  $\nabla$ . From here, the direct sum decomposition can be written as

$$TTM = VTM \oplus HTM.$$

In each local chart, with putting  $\{E_{(j)}\} = \partial_j - u^s \Gamma_{sj}^h \partial_{\bar{h}}$ ,  $\{E_{(\bar{j})}\} = \partial_{\bar{j}}$ , we introduce a useful frame  $\{E_\beta\} = \{E_{(j)}, E_{(\bar{j})}\}$  adapted to the distributions. If  $X = X^i \frac{\partial}{\partial x^i}$  is the local expressions in  $U$  of a vector field  $X$  on  $M$ , then the vertical lift  ${}^V X$ , the horizontal lift  ${}^H X$  and the complete lift  ${}^C X$  of  $X$  are given with respect to the adapted frame  $\{E_\beta\}$  by

$${}^H X = \begin{pmatrix} {}^H X^j \\ {}^H X^{\bar{j}} \end{pmatrix} = \begin{pmatrix} X^j \\ 0 \end{pmatrix}, \tag{1}$$

$${}^C X = \begin{pmatrix} {}^C X^j \\ {}^C X^{\bar{j}} \end{pmatrix} = \begin{pmatrix} X^j \\ u^s \nabla_s X^j \end{pmatrix}. \tag{2}$$

**Lemma 1.** *The Lie brackets of the adapted frame of  $TM$  satisfy the following identities:*

$$\begin{cases} [E_j, E_i] = u^s R_{ijs}{}^m E_{\bar{h}}, \\ [E_j, E_{\bar{i}}] = \Gamma_{ji}^h E_{\bar{h}}, \\ [E_{\bar{j}}, E_{\bar{i}}] = 0. \end{cases}$$

where  $R_{ijs}{}^m$  denotes the components of the curvature tensor of  $M$ .

### 3 Kaluza-Klein metric on tangent bundle

**Definition 1.** *Let  $(M, g)$  be a Riemannian manifold, a metric  ${}^{KK}g$  on  $TM$  will be called Kaluza-Klein if it takes the form*

$${}^{KK}g({}^H X, {}^H Y) = c(t)g(X, Y),$$

$${}^{KK}g({}^H X, {}^V Y) = 0,$$

$${}^{KK}g({}^V X, {}^V Y) = a(t)g(X, Y) + b(t)g(X, Y)$$

where  $t$  is energy density in direction of  $u$  is defined by  $t = g(u, u)/2$ ,  $c$  is strictly positive and  $a$  and  $b$  such that  ${}^{KK}g$  is positive definite.

The metric  ${}^{KK}g$  and its inverse have respectively the components with respect to adapted frame  $E_\beta$  as follows

$${}^{KK}g = \begin{pmatrix} cg_{ji} & 0 \\ 0 & ag_{ji} + bg_{js}g_{it}u^s u^t \end{pmatrix},$$

$$({}^{KK}g)^{-1} = \begin{pmatrix} \frac{1}{c}g^{ij} & 0 \\ 0 & \frac{1}{a}g^{ij} - \frac{b}{a(a+2bt)}u^i u^j \end{pmatrix}.$$

Standard computations with the Koszul's formula and having mind the Lemma 1 give the following proposition.

**Proposition 1.** Let  ${}^{KK}g$  be a Kaluza-Klein metric on  $TM$ , then the corresponding Levi-Civita connection coefficients are given by

$$\begin{cases} {}^{KK}\Gamma_{ji}^h = \Gamma_{ji}^h, & {}^{KK}\Gamma_{ji}^{\bar{h}} = -\frac{1}{2}u^k R_{jik}^h - \frac{c'}{2(a+2tb)}g_{ji}u^h, \\ {}^{KK}\Gamma_{ji}^h = \frac{a}{2c}u^k R_{kij}^h + \frac{c'}{2c}u_i \delta_j^h, & {}^{KK}\Gamma_{ji}^{\bar{h}} = \Gamma_{ji}^h, \\ {}^{KK}\Gamma_{ji}^h = \frac{a}{2}u^k R_{kji}^h + \frac{c}{2c}u_j \delta_i^h, & {}^{KK}\Gamma_{ji}^{\bar{h}} = 0, \\ {}^{KK}\Gamma_{ji}^h = 0, \\ {}^{KK}\Gamma_{ji}^{\bar{h}} = L(u_j \delta_i^h + u_i \delta_j^h) + M g_{ji} u^h + N u_j u_i u^h \end{cases} \quad (3)$$

with respect to adapted frame  $E_\beta$ . Here,  $\Gamma_{ji}^h$  and  $R_{kij}^h$  denotes the Levi-Civita connection and Riemannian curvature components of  $g$  respectively and  $L = \frac{a'}{2a}$ ,  $M = \frac{2b-a'}{2(a+2tb)}$ ,  $N = \frac{ab'-2a'b}{2a(a+2tb)}$ .

For the invariant expression of the Levi-Civita connection, see [6].

#### 4 Some special vector fields on $(TM, {}^{KK}g)$

First, let us compute the covariant derivatives of horizontal lift and complete lift of the vector field  $X$ . Then we have

$${}^{KK}\nabla^H X = \begin{pmatrix} \nabla_i X^h & -\frac{a}{2}R_{iks}^h u^k X^s + \frac{c'}{2c}u_i X^h \\ -\frac{1}{2}R_{isk}^h u^k X^s - \frac{c'}{2(a+2tb)}u^h X_i & 0 \end{pmatrix}, \quad (4)$$

$${}^{KK}\nabla^C X = \begin{pmatrix} \nabla_i X^h - \frac{a}{2c}R_{mki}^h u^k u^s \nabla_s X^m + \frac{c'}{2c}u_m u^s \nabla_s X^m \delta_i^h & -\frac{a}{2}R_{ikm}^h u^k X^m + \frac{c'}{2c}u_i X^h \\ u^s (\nabla_i \nabla_s X^h - \frac{1}{2}R_{ims}^h X^h) - \frac{c'}{2(a+2tb)}u^h X_i & \nabla_i X^h + (L(u_i \delta_m^h + u_m \delta_i^h) + M g_{mi} u^h + N u_m u_i u^h) u^s \nabla_s X^m \end{pmatrix}. \quad (5)$$

Since  $R_{mki}^h X^i = 0$  is a consequence of the fact that  $\nabla_i X^h = 0$ , we obtain the following proposition from (4) and (5).

**Proposition 2.** The complete (resp. horizontal) lift of a vector field on  $M$  to  $TM$  is parallel in  $(TM, {}^{KK}g)$  if the given vector field is parallel and  $c$  is a constant function.

Secondly, we consider the divergence and rotation of the lifts of vector fields. A vector field  $X$  on  $(M, g)$  is said to be *incompressible* if and only if its divergence  $div X = \nabla_j X^j = 0$ . On the other hand, a vector field  $X$  on  $(M, g)$  is said to be *closed* if the rotation of associated covector field (whose components are  $X_j = g_{ji} X^i$ ) is zero, i.e.,

$$(rot X)_{ji} = \partial_j X_i - \partial_i X_j = \nabla_j X_i - \nabla_i X_j = 0. \quad (6)$$

$X$  is said to be *harmonic* on  $(M, g)$  if it is both incompressible and closed. From (4) and (5), the divergence of the lifts are

$$div^H X = div X, \quad div^C X = 2div X + u_m [L(n+1) + M + N \|u\|^2] u^s \nabla_s X^m + \frac{nc'}{2c} u_m u^s \nabla_s X^m. \quad (7)$$

Thus, it is clear that the horizontal lift of  $X$  is incompressible if and only if it is incompressible in  $M$ . But, the complete lift of  $X$  is not incompressible even if it is incompressible in  $M$  in general. For the special case  $a = cons.$ ,  $b = 0$  and  $c = cons.$ , we can say the complete lift of  $X$  is incompressible if and only if it is incompressible in  $M$ .

The rotation of associated covector field of a vector field  $\tilde{X}$  on  $(TM, {}^{KK}\nabla)$  is given by

$${}^{KK}\nabla_j X_i - {}^{KK}\nabla_i X_j = {}^{KK}g_{IM} {}^{KK}\nabla_j X^M - {}^{KK}g_{JM} {}^{KK}\nabla_i X^M. \quad (8)$$

Using (4) and (5), we obtain the components of  $rot^H X$  and  $rot^C X$  as

$$\begin{aligned}
 rot^H X_{ji} &= c(rot X)_{ji}, \\
 rot^H X_{\bar{j}i} &= \left(\frac{a(1-c)}{2}R_{jki} + \frac{b}{2}u_j u_s R_{ikb}^s\right)u^k X^b + \frac{c'}{2(a+2tb)}(a+b\|u\|^2)u_j X_i, \\
 rot^H X_{\bar{j}\bar{i}} &= \left(\frac{a(1-c)}{2}R_{bjki} + \frac{b}{2}u_i u_s R_{bjk}^s\right)u^k X^b - \frac{c'}{2(a+2tb)}(a+b\|u\|^2)u_j X_i, \\
 rot^H X_{\bar{j}\bar{i}} &= 0; \\
 rot^C X_{ji} &= c(rot X)_{ji} + aR_{jikm}u^k u^b (\nabla_b X^m), \\
 rot^C X_{\bar{j}i} &= \left(\frac{a(1-c)}{2}R_{ikbj} + \frac{b}{2}u_j u_s R_{ikb}^s\right)u^k X^b - u^s (a\nabla_i \nabla_s X_j + b\nabla_i \nabla_s X^a u_j u_a) + \frac{c'}{2(a+2tb)}(a+b\|u\|^2)u_j X_i, \\
 rot^C X_{\bar{j}\bar{i}} &= \left(\frac{a(1-c)}{2}R_{bjki} + \frac{b}{2}u_i u_s R_{bjk}^s\right)u^k X^b + u^s (a\nabla_j \nabla_s X_i + b\nabla_j \nabla_s X^a u_i u_a) - \frac{c'}{2(a+2tb)}(a+b\|u\|^2)u_i X_j, \\
 rot^C X_{\bar{j}\bar{i}} &= a(rot X)_{ji} + bu^s (\nabla_j X_s u_i - \nabla_i X_s u_j).
 \end{aligned} \tag{9}$$

Thus, we get the following proposition.

**Proposition 3.** *Let  $(M, g)$  be a flat Riemannian manifold and  $c$  is a constant function. Then the followings hold.*

- (i) *The horizontal lift of a vector field  $X$  on  $M$  to  $(TM,^{KK}g)$  is harmonic if and only if  $X$  is harmonic on  $M$ .*
- (ii) *The complete lift of a vector field  $X$  is harmonic if and only if it is parallel in  $(M, g)$ .*

### 5 Geodesics on $(TM,^{KK}g)$

It is known that a curve  $\tilde{\gamma}$  is a geodesic on  $TM$  with respect to  $^{KK}\nabla$  if and only if it satisfies the following differential equations

$$\frac{d}{dt}\left(\frac{\theta^\alpha}{dt}\right) + ^{KK}\Gamma_{\gamma\beta}^\alpha \frac{\theta^\gamma}{dt} \frac{\theta^\beta}{dt} = 0$$

with respect to adapted frame, where

$$\frac{\theta^h}{dt} = \frac{dx^h}{dt}, \quad \frac{\theta^{\bar{h}}}{dt} = \frac{\delta u^h}{dt} = \frac{du^h}{dt} + \Gamma_{ji}^h \frac{dx^j}{dt} u^i.$$

By direct computations with using (4), then we have the following proposition.

**Proposition 4.** *Let  $\tilde{\gamma}$  be a curve on  $TM$  with locally expressions  $x^h = x^h(t), x^{\bar{h}} = u^h(t)$  with respect to induced coordinates  $(x^h, x^{\bar{h}})$  in  $\pi^{-1}(U) \subset TM$ . Then the curve  $\tilde{\gamma}$  is a geodesic on  $(TM,^{KK}g)$  if and only if it satisfies the following equations*

$$\begin{cases} \frac{\delta^2 x^h}{dt^2} + \frac{a(c+1)}{2c}(u^k R_{kji}^h + \frac{c'}{c}u_j \delta_i^h) \frac{\delta u^j}{dt} \frac{dx^i}{dt} = 0, \\ \frac{\delta^2 u^h}{dt^2} + [L(u_j \delta_i^h + u_i \delta_j^h) + Mg_{ji}u^h + Nu_j u_i u^h] \frac{\delta u^j}{dt} \frac{\delta u^i}{dt} - \frac{c'}{2(a+2tb)}g_{ji}u^h \frac{dx^j}{dt} \frac{dx^i}{dt} = 0. \end{cases} \tag{11}$$

If a curve  $\tilde{\gamma}$  satisfying the equations (11) lies on a fibre given by  $x^h = cons.$ , then by virtue of  $\frac{dx^h}{dt} = 0$  and  $\frac{\delta u^h}{dt} = \frac{du^h}{dt}$ , the equations (11) reduces to

$$\frac{\delta^2 u^h}{dt^2} + [L(u_j \delta_i^h + u_i \delta_j^h) + Mg_{ji}u^h + Nu_j u_i u^h] \frac{du^j}{dt} \frac{du^i}{dt} = 0. \tag{12}$$

Hence, we get the proposition below.

**Proposition 5.** *If a geodesic lies on a fibre of  $(TM,^{KK}g)$ , it's local expression as in (12).*

Now, let  $\gamma = \pi \circ^H \gamma$  be a geodesic on  $(M, \nabla)$ , i.e.  $\frac{\delta^2 x^h}{dt^2} = 0$ . Using this condition and  $\frac{\delta u^j}{dt} = \frac{\delta x^h}{dt} = 0$ , the following proposition is obtained.

**Proposition 6.** *The horizontal lift of a geodesic on  $(M, \nabla)$  is not a geodesic  $(TM, {}^{KK}g)$ .*

Finally, let  $\gamma = \pi \circ \tilde{\gamma}$  be a geodesic on  $(M, \nabla)$ , i.e.  $\frac{\delta^2 x^h}{dt^2} = \frac{\delta}{dt}(\frac{dx^h}{dt}) = 0$ . On the other hand, from the definition of the natural lift of the curve  $(x^h = x^h(t), u^h = \frac{dx^h}{dt})$ , we obtain

$$\frac{\delta^2 u^h}{dt^2} = \frac{c'}{2(a+2tb)} g_{ji} u^h \frac{dx^j}{dt} \frac{dx^i}{dt} \quad (13)$$

By the virtue of (13) and (11), it is easily seen that the natural lift of a curve on  $M$  is not a geodesic on  $(TM, {}^{KK}g)$ . This corollary is given as a proposition below.

**Proposition 7.** *The natural lift  $\tilde{\gamma}$  any geodesic on  $M$  is not a geodesic on  $(TM, {}^{KK}g)$ .*

## 6 Conclusions

In this study, some characterizations on specific vector fields and geodesics are given on tangent bundle with respect to a Cheeger-Gromoll type metric, say Kaluza-Klein metric. The results in this paper have two significant difference from previous papers. Firstly, the proposition 3 was given as follows with respect to Sasaki metric in [7] : "The complete (resp. horizontal) lift of a vector field  $X$  on  $M$  to  $TM$  is harmonic with in  $(TM, {}^Sg)$  if and only if  $X$  is harmonic and has vanishing second covariant derivative (resp. harmonic) in  $(M, g)$ ." Here, this proposition is written under the additional conditions such that flatness of the base manifold and the constancy of the function  $c$ . Secondly, the last two propositions on geodesics were written affirmatively with respect to the metrics Sasaki [8], Cheeger-Gromoll [9] and the metric  $g_{a,b}$  [10].

## Competing interests

The authors declare that they have no competing interests.

## Authors' contributions

All authors have contributed to all parts of the article. All authors read and approved the final manuscript.

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