# Some characterizations of constant breadth timelike curves in Minkowski 4-space $E_{1}^{4}$ 

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#### Abstract

In this study, the differential equation characterizations of constant breadth timelike curves are given in Minkowski 4-space $E_{1}^{4}$. Furthermore, a criterion for a timelike curve to be a curve of constant breadth in $E_{1}^{4}$ is introduced. As an example, the obtained results are applied to the case that the curvatures $k_{1}, k_{2}, k_{3}$ and are discussed.


Keywords: Constant breadth curve, timelike curve, Frenet frame.

## 1 Introduction

Euler introduced the constant breadth curves in 1778 [6]. He considered these special curves in the plane. Later, many geometers have shown increased interest in the properties of plane convex curves. Struik published a brief review of the most important publications on this subject [19]. Also, Ball [1], Barbier [2], Blaschke [3,4] and Mellish [13] investigated the properties of plane curves of constant breadth. A space curve of constant breadth was obtained by Fujiwara by taking a closed curve whose normal plane at a point $P$ has only one more point $Q$ in common with the curve, and for which the distance $d(P, Q)$ is constant [7]. He also defined and studied constant breadth surfaces. Later, Smakal studied the constant breadth space curves [18]. Furthermore, Blaschke considered the notion of curve of constant breadth on the sphere [4]. Moreover, Reuleaux studied the curves of constant breadth and gave the method related to these curves for the kinematics of machinery [15]. Then, constant breadth curves had an importance for engineering sciences and by considering this fact Tanaka used the constant breadth curves in the kinematics design of Com follower systems [20].

Moreover, Köse has presented some concepts for space curves of constant breadth in Euclidean 3-space in [10] and Sezer has obtained the differential equations characterizing space curves of constant breadth and introduced a criterion for these curves [17]. Constant breadth curves in Euclidean 4-space were given by Mağden and Köse [11]. Moreover, constant breath curves have been studied in Minkowski space. Kazaz, Önder and Kocayiğit have studied spacelike curves of constant breadth in Minkowski 4-space [8]. Later, Önder, Kocayiğit and Candan have obtained and studied the differential equations characterizing constant breadth curves in Minkowski 3-space [14]. Furthermore, Kocayiğit and Önder have showed that constant breadth spacelike curves are normal curves, helices, and spherical curves in some special cases im Minkowski 3-space [9]. Moreover, in [12] Mağden and Yılmaz have given characterizations curves of constant breadth in four dimensional Galilean space in terms Frenet-Serret vector fields. Also, Yılmaz and Turgut have presented partially null curves of constant breadth in Semi-Riemannian space [21].

In this paper, we study the differential equations characterizing constant breadth timelike curves in the Minkowski 4 -space $E_{1}^{4}$. Moreover, we give a criterion characterizing these curves in $E_{1}^{4}$.

## 2 Differential equations characterizing constant breadth timelike curves in $E_{1}^{4}$

Let $(C)$ be a unit speed regular timelike curve in the Minkowski 4-space $E_{1}^{4}$ with parametrization $\alpha(s): I \subset \mathbb{R} \rightarrow E_{1}^{4}$. Denote by $\{\mathbf{T}, \mathbf{N}, \mathbf{B}, \mathbf{E}\}$ the moving Frenet frame along the timelike curve $(C)$ in $E_{1}^{4}$. Then, the following Frenet formulae are given,

$$
\left[\begin{array}{c}
\mathbf{T}^{\prime} \\
\mathbf{N}^{\prime} \\
\mathbf{B}^{\prime} \\
\mathbf{E}^{\prime}
\end{array}\right]=\left[\begin{array}{cccc}
0 & k_{1} & 0 & 0 \\
k_{1} & 0 & k_{2} & 0 \\
0 & -k_{2} & 0 & k_{3} \\
0 & 0 & -k_{3} & 0
\end{array}\right]\left[\begin{array}{c}
\mathbf{T} \\
\mathbf{N} \\
\mathbf{B} \\
\mathbf{E}
\end{array}\right]
$$

where $k_{1}, k_{2}$ and $k_{3}$ are the first, second and third curvatures of the curve $(C)$, respectively and $\{\mathbf{T}, \mathbf{N}, \mathbf{B}, \mathbf{E}\}$ denote the tangent, the principal normal, the first binormal and the second binormal vector fields, respectively and they satisfy the following equalities:

$$
<\mathbf{T}, \mathbf{T}>=-1,<\mathbf{N}, \mathbf{N}>=<\mathbf{B}, \mathbf{B}>=<\mathbf{E}, \mathbf{E}>=1
$$

where $<,>$ is the Lorentzian inner product defined by

$$
<\mathbf{a}, \mathbf{b}>=-a_{1} b_{1}+a_{2} b_{2}+a_{3} b_{3}+a_{4} b_{4}
$$

here $\mathbf{a}=\left(a_{1}, a_{2}, a_{3}, a_{4}\right), \mathbf{b}=\left(b_{1}, b_{2}, b_{3}, b_{4}\right)$ are the vectors in $E_{1}^{4}[22]$.

Definition 1. Let $(C)$ be a unit speed regular timelike curve in $E_{1}^{4}$ with position vector $\alpha(s)$. If $(C)$ has parallel tangents $\mathbf{T}$ and $\mathbf{T}^{*}$ in opposite direction at the opposite points $\alpha$ and $\alpha^{*}$ of the curve and if the distance between these points is always constant then $(C)$ is called a timelike curve of constant breadth in $E_{1}^{4}$. Moreover, a pair of curves $(C)$ and $\left(C^{*}\right)$ for which the tangents at the corresponding points are parallel and in opposite directions and the distance between these points is always constant is called a timelike curve pair of constant breadth in $E_{1}^{4}$.

Let now $(C)$ and $\left(C^{*}\right)$ be a pair of unit speed curves in $E_{1}^{4}$ with position vectors $\alpha(s)$ and $\alpha^{*}\left(s^{*}\right)$, where $s$ and $s^{*}$ are arc length parameters of the curves, respectively. Let $(C)$ and $\left(C^{*}\right)$ have parallel tangents in opposite directions at opposite points. Then the curve $\left(C^{*}\right)$ may be represented by the equation

$$
\begin{equation*}
\alpha^{*}(s)=\alpha(s)+m_{1}(s) \mathbf{T}(s)+m_{2}(s) \mathbf{N}(s)+m_{3}(s) \mathbf{B}(s)+m_{4}(s) \mathbf{E}(s) \tag{1}
\end{equation*}
$$

where $m_{i}(s),(1 \leq i \leq 4)$ are the differentiable functions of $s$ which is the arc length of $(C)$. Differentiating this equation with respect to $s$ and using the Frenet formulae we obtain

$$
\begin{aligned}
\frac{\alpha^{*}(s)}{d s}=\mathbf{T}^{*} \frac{d s^{*}}{d s} & =\left(1+\frac{d m_{1}}{d s}+m_{2} k_{1}\right) \mathbf{T}+\left(m_{1} k_{1}+\frac{d m_{2}}{d s}-m_{3} k_{2}\right) \mathbf{N} \\
& +\left(m_{2} k_{2}+\frac{d m_{3}}{d s}-m_{4} k_{3}\right) \mathbf{B}+\left(m_{3} k_{3}+\frac{d m_{4}}{d s}\right) \mathbf{E}
\end{aligned}
$$

Since $\mathbf{T}=-\mathbf{T}^{*}$ at the corresponding points of $(C)$ and $\left(C^{*}\right)$, we have

$$
\left\{\begin{array}{l}
\left(1+\frac{d m_{1}}{d s}+m_{2} k_{1}\right)=-\frac{d s^{*}}{d s}  \tag{2}\\
\left(m_{1} k_{1}+\frac{d m_{2}}{d s}-m_{3} k_{2}\right)=0, \\
\left(m_{2} k_{2}+\frac{d m_{3}}{d s}-m_{4} k_{3}\right)=0, \\
\left(m_{3} k_{3}+\frac{d m_{4}}{d s}\right)=0
\end{array}\right.
$$

It is well known that the curvature of $(C)$ is $\lim (\Delta \varphi / \Delta s)=(d \varphi / d s)=k_{1}(s)$, where $\varphi=\int_{0}^{s} k_{1}(s) d s$ is the angle between the tangent of the curve $(C)$ and a given fixed direction at the point $\alpha(s)$. Then from (2) we have the following system

$$
\begin{equation*}
m_{1}^{\prime}=-m_{2}-f(\varphi), m_{2}^{\prime}=m_{3} \rho k_{2}-m_{1}, m_{3}^{\prime}=m_{4} \rho k_{3}-m_{2} \rho k_{2}, m_{4}^{\prime}=-m_{3} \rho k_{3} . \tag{3}
\end{equation*}
$$

Here and after we will use $\left({ }^{\prime}\right)$ to show the differentiation with respect to $\varphi \cdot \operatorname{In}(3), f(\varphi)=\rho+\rho^{*}$ and, $\rho=\frac{1}{k_{1}}$ and $\rho^{*}=\frac{1}{k_{1}^{*}}$ denote the radius of curvatures at the points $\alpha$ and $\alpha^{*}$, respectively. From (3) eliminating $m_{2}, m_{3}$ and $m_{4}$ their derivatives we have the following differential equation

$$
\begin{equation*}
\frac{d}{d \varphi}\left[\frac{1}{\rho k_{3}} \frac{d}{d \varphi}\left[\frac{1}{\rho k_{2}}\left(\frac{d^{2} m_{1}}{d \varphi^{2}}-m_{1}\right)\right]-\frac{k_{2}}{k_{3}} \frac{d m_{1}}{d \varphi}\right]+\frac{k_{3}}{k_{2}}\left(\frac{d^{2} m_{1}}{d \varphi^{2}}-m_{1}\right)+\frac{d}{d \varphi}\left[\frac{1}{\rho k_{3}} \frac{d}{d \varphi}\left(\frac{1}{\rho k_{2}} \frac{d f}{d \varphi}\right)-\frac{k_{2}}{k_{3}} f\right]+\frac{k_{3}}{k_{2}} \frac{d f}{d \varphi}=0 \tag{4}
\end{equation*}
$$

Then we can give the following theorem.

Theorem 1. The general differential equation characterizing constant breadth timelike curves in $E_{1}^{4}$ is given by (4).

Let now consider the system (3) again. The distance $d$ between the opposite points $\alpha$ and $\alpha^{*}$ is the breadth of the curves and is constant, that is,

$$
\begin{equation*}
d^{2}=\|\mathbf{d}\|^{2}=\left\|\alpha^{*}-\alpha\right\|^{2}=-m_{1}^{2}+m_{2}^{2}+m_{3}^{2}+m_{4}^{2}=\text { const } . \tag{5}
\end{equation*}
$$

Then the system (3) may be written as follows:

$$
\begin{equation*}
m_{2}=-f(\varphi), m_{2}^{\prime}=m_{3} \rho k_{2}, m_{3}^{\prime}=m_{4} \rho k_{3}-m_{2} \rho k_{2}, m_{4}^{\prime}=-m_{3} \rho k_{3}, m_{1}=0 \tag{6}
\end{equation*}
$$

or

$$
\begin{equation*}
m_{1}^{\prime}=-m_{2}, m_{2}^{\prime}=-m_{1}+m_{3} \rho k_{2}, m_{3}^{\prime}=m_{4} \rho k_{3}-m_{2} \rho k_{2}, m_{4}^{\prime}=-m_{3} \rho k_{3} \tag{7}
\end{equation*}
$$

which are the systems describing the curve (1).

Let us consider the system (7) with special chosen $m_{1}=$ const. . Here, eliminating first $m_{1}, m_{2}, m_{3}$ and their derivatives, and then $m_{1}, m_{2}, m_{4}$ and their derivatives, respectively, we obtain the following linear differential equations of second order

$$
\begin{cases}\left(\rho k_{3}\right) m_{4}^{\prime \prime}-\left(\rho k_{3}\right)^{\prime} m_{4}^{\prime}+\left(\rho k_{3}\right)^{3} m_{4}=0, & \rho k_{2} \neq 0  \tag{8}\\ \left(\rho k_{3}\right) m_{3}^{\prime \prime}-\left(\rho k_{3}\right)^{\prime} m_{3}^{\prime}+\left(\rho k_{3}\right)^{3} m_{3}=0, & \rho k_{3} \neq 0\end{cases}
$$

By changing the variable $\varphi$ of the form $\xi=\int_{0}^{\varphi} \rho(t) k_{3}(t) d t$, these equations can be transformed into the following differential equations with constant coefficients,

$$
\begin{equation*}
\frac{d^{2} m_{4}}{d \xi^{2}}+m_{4}=0 \text { and } \frac{d^{2} m_{3}}{d \xi^{2}}+m_{3}=0 \tag{9}
\end{equation*}
$$

respectively. Then, the general solutions of the differential equations in (9) are

$$
\left\{\begin{array}{l}
m_{3}=A \cos \left(\int_{0}^{\varphi} \rho k_{3} d t\right)+B \sin \left(\int_{0}^{\varphi} \rho k_{3} d t\right)  \tag{10}\\
m_{4}=C \cos \left(\int_{0}^{\varphi} \rho k_{3} d t\right)+D \sin \left(\int_{0}^{\varphi} \rho k_{3} d t\right)
\end{array}\right.
$$

respectively, where $A, B, C$ and $D$ are real constants. Substituting (10) into (7), we obtain $A=-D, B=C$, and so, the set of the solutions of the system (7), in the form

$$
\left\{\begin{array}{c}
m_{1}=c=\text { const. }, m_{2}=0,  \tag{11}\\
m_{3}=A \cos \int_{0}^{\varphi} \rho k_{3} d t+B \sin \int_{0}^{\varphi} \rho k_{3} d t \\
m_{4}=B \cos \int_{0}^{\varphi} \rho k_{3} d t-A \sin \int_{0}^{\varphi} \rho k_{3} d t
\end{array}\right\}
$$

Thus the equation (1) is described and since $d^{2}=\left\|\alpha^{*}-\alpha\right\|^{2}=$ const., from (11) the breadth of the curve is $d^{2}=-c^{2}+A^{2}+B^{2}$.

Now, let us return to the system (6) with $m_{1}=0$. By changing the variable $\varphi$ of the form $u=\int_{0}^{\varphi} \mu(t) d t, \mu=\rho k_{3}$ and eliminating $m_{1}, m_{2}, m_{4}$ and their derivatives we have the linear differential equation

$$
\begin{equation*}
\frac{d^{2} m_{3}}{d u^{2}}+m_{3}=-\frac{d}{d u}\left(\frac{k_{2}}{k_{3}} m_{2}\right) \tag{12}
\end{equation*}
$$

which has the following solution

$$
\begin{equation*}
m_{3}=A_{1} \cos \int_{0}^{\varphi} \rho k_{3} d t+B_{1} \sin \int_{0}^{\varphi} \rho k_{3} d t-\int_{0}^{\varphi} \cos [u(\varphi)-u(t)] \rho k_{2} f(t) d t \tag{13}
\end{equation*}
$$

Then, the general solution of the system (6) is

$$
\left\{\begin{array}{l}
m_{1}=0  \tag{14}\\
m_{2}=f(\varphi) \\
m_{3}=A_{1} \cos \int_{0}^{\varphi} \rho k_{3} d t+B_{1} \sin \int_{0}^{\varphi} \rho k_{3} d t+\int_{0}^{\varphi} \cos [u(\varphi)-u(t)] \rho k_{2} f(t) d t \\
m_{4}=B_{1} \cos \int_{0}^{\varphi} \rho k_{3} d t-A_{1} \sin \int_{0}^{\varphi} \rho k_{3} d t-\int_{0}^{\varphi} \sin [u(\varphi)-u(t)] \rho k_{2} f(t) d t
\end{array}\right.
$$

which determines the constant breadth timelike curve in (1) where $A_{1}, B_{1}$ are real constants.
Furthermore, in this case, i.e., $m_{1}=0$, from (4) we have the following differential equation

$$
\begin{equation*}
\frac{d}{d \varphi}\left[\frac{1}{\rho k_{3}} \frac{d}{d \varphi}\left(\frac{1}{\rho k_{2}} \frac{d f}{d \varphi}\right)-\frac{k_{2}}{k_{3}} f\right]+\frac{k_{2}}{k_{3}} \frac{d f}{d \varphi}=0 \tag{15}
\end{equation*}
$$

By changing the variable $\varphi$ of the form $w=\int_{0}^{\varphi} \rho k_{2} d \varphi$, (15) becomes

$$
\begin{equation*}
\frac{d}{d w}\left[\frac{k_{2}}{k_{3}}\left(\frac{d^{2} f}{d w^{2}}-f\right)\right]+\frac{k_{3}}{k_{2}} \frac{d f}{d w}=0 \tag{16}
\end{equation*}
$$

which also determines the constant breadth curve in (1).

So far we have dealt with a pair of timelike space curves having parallel tangents in opposite directions at corresponding
points. Now let us consider a simple closed unit speed timelike space curve $(C)$ in $E_{1}^{4}$ for which the normal plane of every point $P$ on the curve meets the curve of a single opposite point $Q$ other than $P$. Then, we may give the following theorem concerning the constant breadth timelike space curves in $E_{1}^{4}$.

Theorem 2. Let (C) be a closed timelike space curve in $E_{1}^{4}$ having parallel tangents in opposite directions at the opposite points of the curve. If the chord joining the opposite points of $(C)$ is a double-normal if and only if $(C)$ is a timelike curve of constant breadth in $E_{1}^{4}$.

Proof. Let the vector $\mathbf{d}=\alpha^{*}-\alpha=m_{1} \mathbf{T}+m_{2} \mathbf{N}+m_{3} \mathbf{B}+m_{4} \mathbf{E}$ be a double-normal of $(C)$ where $m_{1}, m_{2}, m_{3}$ and $m_{4}$ are the functions of $s$, the arc length parameter of the curve. Then we get $\left\langle\mathbf{d}, \mathbf{T}^{*}\right\rangle=-\langle\mathbf{d}, \mathbf{T}\rangle=m_{1}=0$. Thus from (2) we have

$$
\begin{equation*}
m_{2} \frac{d m_{2}}{d s}+m_{3} \frac{d m_{3}}{d s}+m_{4} \frac{d m_{4}}{d s}=0 \tag{17}
\end{equation*}
$$

It follows that $m_{2}^{2}+m_{3}^{2}+m_{4}^{2}=$ constant, i.e., the breadth of $(C)$ is constant, i.e., $(C)$ is a constant breadth timelike curve in $E_{1}^{4}$.

Conversely, if $(C)$ is a constant breadth timelike curve in $E_{1}^{4}$ then $\|\mathbf{d}\|^{2}=-m_{1}^{2}+m_{2}^{2}+m_{3}^{2}+m_{4}^{2}=$ constant. Then as shown, $m_{1}=0$. This means that $\mathbf{d}$ is perpendicular to both $\mathbf{T}$ and $\mathbf{T}^{*}$. So, $\mathbf{d}$ is the double-normal of $(C)$.

A simple closed timelike curve having parallel tangents in opposite directions at opposite points may be represented by the system (14). In this case a pair of opposite points of the curve is $\left(\alpha^{*}(\varphi), \alpha(\varphi)\right)$ for $\varphi$, where $0 \leq \varphi \leq 2 \pi$. Since ( $C$ ) is a simple closed timelike curve we get $\alpha^{*}(0)=\alpha^{*}(2 \pi)$. Hence from (14) we have

$$
\begin{equation*}
\int_{0}^{2 \pi} \rho k_{3} d t=2 n \pi, \quad(n \in \mathbb{Z}) \tag{18}
\end{equation*}
$$

Using the equality $d s=\rho d \varphi$, this formula may be given as $\int_{C} k_{3} d s=2 n \pi,(n \in \mathbb{Z})$. This says that the integral third curvature of $(C)$ is zero. So, we can give the following corollary.

Corollary 1. The total third curvature of a simple closed timelike curve (C) of constant breadth is $2 n \pi, n \in \mathbb{Z}$.
Furthermore, if we take $\frac{k_{2}}{k_{3}}=a=$ constant, then from (16) we have

$$
\begin{equation*}
\frac{d^{3} f}{d w^{3}}+K \frac{d f}{d w}=0 \tag{19}
\end{equation*}
$$

where $K=-1+\frac{1}{a^{2}}$. If we assume $K \neq \pm 1$, the general solution of (19) is

$$
\begin{equation*}
f=A_{2} \sin \int_{0}^{\varphi} K \rho k_{2} d t+B_{2} \cos \int_{0}^{\varphi} K \rho k_{2} d t+C_{1} \tag{20}
\end{equation*}
$$

where $A_{2}, B_{2}$ and $C_{1}$ are real constants. Since $(C)$ is a simple closed timelike curve, i.e., $\alpha^{*}(0)=\alpha^{*}(2 \pi)$, from (20) it follows,

$$
\begin{equation*}
\int_{0}^{\varphi} K \rho k_{2} d t=2 n \pi, \quad(n \in \mathbb{Z}) \tag{21}
\end{equation*}
$$

Using the equality $d s=\rho d \varphi$, this formula may be given as $\int_{C} k_{2} d s=2 \frac{n}{K} \pi, \quad(K, n \in \mathbb{Z})$. This says that the integral second curvature of $(C)$ is $2 \frac{n}{K} \pi, \quad(K, n \in \mathbb{Z})$. So, we can give the following corollary.

Corollary 2. The total second curvature of a simple closed constant breadth timelike curve ( $C$ ) with $a=k_{2} / k_{3}=$ constant is $2 \frac{n}{K} \pi$, where $n \in \mathbb{Z}$ and $K=-1+\frac{1}{a^{2}}$.

## 3 A criterion for constant breadth timelike curves in $E_{1}^{4}$

Let us assume that $(C)$ is a constant breadth timelike curve in $E_{1}^{4}$ and $\alpha(s)$ denotes the position vector of a generic point of the curve. If $(C)$ is a closed curve, the position vector $\alpha(s)$ must be a periodic function of period $\omega=2 \pi$, where $\omega$ is the total length of $(C)$. Then the curvatures $k_{1}(s), k_{2}(s)$ and $k_{3}(s)$ are also periodic of the same period. However, periodicity of the curvatures and closeness of the curve are not sufficient to guarantee that a timelike space curve is a constant breadth curve in $E_{1}^{4}$. That is, if a timelike curve is closed curve (periodic), it may be the constant breadth curve or not. Therefore, to guarantee that a timelike curve is a constant breadth curve, we may use the system (7) characterizing a constant breadth timelike curve and follow the similar way given in [5].

For this purpose, first let us consider the following Frenet formulas at a generic point on the curve $(C)$,

$$
\begin{equation*}
\frac{d \mathbf{T}}{d s}=k_{1} \mathbf{N}, \frac{d \mathbf{N}}{d s}=k_{1} \mathbf{T}+k_{2} \mathbf{B}, \frac{d \mathbf{B}}{d s}=-k_{2} \mathbf{N}+k_{3} \mathbf{E}, \frac{d \mathbf{E}}{d s}=-k_{3} \mathbf{B} . \tag{22}
\end{equation*}
$$

Writing the formulas (22) in terms of $\varphi$ and allowing for $\frac{d \varphi}{d s}=k_{1}=\frac{1}{\rho}$ we have

$$
\begin{equation*}
\frac{d \mathbf{T}}{d \varphi}=\mathbf{N}, \frac{d \mathbf{N}}{d \varphi}=\mathbf{T}+\rho k_{2} \mathbf{B}, \frac{d \mathbf{B}}{d \varphi}=-\rho k_{2} \mathbf{N}+\rho k_{3} \mathbf{E}, \frac{d \mathbf{E}}{d \varphi}=-\rho k_{3} \mathbf{B} . \tag{23}
\end{equation*}
$$

Furthermore we can write the Frenet vectors $\mathbf{T}, \mathbf{N}, \mathbf{B}, \mathbf{E}$ in the coordinate forms as follows

$$
\begin{equation*}
\mathbf{T}=\sum_{i=1}^{4} t_{i} \mathbf{e}_{\mathbf{i}}, \mathbf{N}=\sum_{i=1}^{4} n_{i} \mathbf{e}_{\mathbf{i}}, \mathbf{B}=\sum_{i=1}^{4} b_{i} \mathbf{e}_{\mathbf{i}}, \mathbf{E}=\sum_{i=1}^{4} \varepsilon_{i} \mathbf{e}_{\mathbf{i}} . \tag{24}
\end{equation*}
$$

Since $\{\mathbf{T}, \mathbf{N}, \mathbf{B}, \mathbf{E}\}$ is the orthonormal base in $E_{1}^{4}$, putting (24) and their derivatives into (23), we have the systems of linear differential equations

$$
\left\{\begin{array}{llll}
\frac{d t_{1}}{d \varphi}=n_{1}, & \frac{d t_{2}}{d \varphi}=n_{2}, & \frac{d t_{3}}{d \varphi}=n_{3}, & \frac{d t_{4}}{d \varphi}=n_{4}  \tag{25}\\
\frac{d n_{1}}{d \varphi}=t_{1}+\rho k_{2} b_{1}, & \frac{d n_{2}}{d \varphi}=t_{2}+\rho k_{2} b_{2}, & \frac{d n_{3}}{d \varphi}=t_{3}+\rho k_{2} b_{3}, & \frac{d n_{4}}{d \varphi}=t_{4}+\rho k_{2} b_{4} \\
\frac{d b_{1}}{d \varphi}=\rho k_{3} \varepsilon_{1}-\rho k_{2} n_{1}, \frac{d b_{2}}{d \varphi}=\rho k_{3} \varepsilon_{2}-\rho k_{2} n_{2}, \frac{d b_{3}}{d \varphi}=\rho k_{3} \varepsilon_{3}-\rho k_{2} n_{3}, & \frac{d b_{4}}{d \varphi}=\rho k_{3} \varepsilon_{4}-\rho k_{2} n_{4} \\
\frac{d \varepsilon_{1}}{d \varphi}=-\rho k_{3} b_{1}, & \frac{d \varepsilon_{2}}{d \varphi}=-\rho k_{3} b_{2}, & \frac{d \varepsilon_{3}}{d \varphi}=-\rho k_{3} b_{3}, & \frac{d \varepsilon_{4}}{d \varphi}=-\rho k_{3} b_{4} .
\end{array}\right\}
$$

From (25), we find that $\left\{t_{1}, n_{1}, b_{1}, \varepsilon_{1}\right\},\left\{t_{2}, n_{2}, b_{2}, \varepsilon_{2}\right\},\left\{t_{3}, n_{3}, b_{3}, \varepsilon_{3}\right\}$ and $\left\{t_{4}, n_{4}, b_{4}, \varepsilon_{4}\right\}$ are four independent solutions of the following system of differential equations:

$$
\begin{equation*}
\frac{d \psi_{1}}{d \varphi}=\psi_{2}, \frac{d \psi_{2}}{d \varphi}=\psi_{1}+\rho k_{2} \psi_{3}, \frac{d \psi_{3}}{d \varphi}=\rho k_{3} \psi_{4}-\rho k_{2} \psi_{2}, \frac{d \psi_{4}}{d \varphi}=-\rho k_{3} \psi_{3} \tag{26}
\end{equation*}
$$

If the curve $(C)$ is the constant breadth timelike curve, then the systems (7) and (26) must be the same system. So, we observe that $\psi_{1}=m_{1}, \psi_{2}=m_{2}, \psi_{3}=m_{3}, \psi_{4}=m_{4}$. For brevity, we can write (7) or (26) in the form

$$
\begin{equation*}
\frac{d \psi}{d \varphi}=A(\varphi) \psi \tag{27}
\end{equation*}
$$

where

$$
\psi=\left[\begin{array}{l}
m_{1} \\
m_{2} \\
m_{3} \\
m_{4}
\end{array}\right], A(\varphi)=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
1 & 0 & \rho k_{2} & 0 \\
0-\rho k_{2} & 0 & \rho k_{3} \\
0 & 0 & -\rho k_{3} & 0
\end{array}\right] .
$$

Obviously, (27) is a special case of the general linear differential equations abbreviated to the form

$$
\left\{\begin{array}{l}
\frac{d \psi}{d t}=A(t) \psi  \tag{28}\\
\varphi=\left[\begin{array}{c}
m_{1} \\
m_{2} \\
\vdots \\
m_{n}
\end{array}\right], A(t)=\left[\begin{array}{cccc}
a_{11} & a_{12} & \cdots & a_{1 n} \\
a_{21} & a_{22} & \cdots & a_{2 n} \\
\vdots & \vdots & \ddots & \vdots \\
a_{n 1} & a_{n 2} & \cdots & a_{n n}
\end{array}\right],(4 \leq n)
\end{array}\right.
$$

where $a_{i j}(t)$ are assumed to be continuous and periodic of period $\omega$ (See [5,16]). Let the initial conditions be $\psi_{i}(0)=x_{i}$, $(i=1,2, \ldots, n)$. Let us take $x=\left[x_{1}, x_{2}, \ldots, x_{n}\right]^{T}$ and

$$
\psi(t, x)=\left[m_{1}(t, x), m_{2}(t, x), \ldots, m_{n}(t, x)\right]^{T}
$$

Then the equation (28) may be written in the form $\frac{d \psi}{d t}=A(t) \psi, \psi(0)=x$ as is well known from [5], the solution $\psi(t, x)$ of this equation is periodic of period $\omega$, if

$$
\int_{0}^{\omega} A(\xi) \psi(\xi, x) d \xi=0
$$

and

$$
\left\{\begin{array}{l}
\psi(t, x)=\{E+M(t)\} x,(E=\text { unit matrix })  \tag{29}\\
M(t)=I A(t)+I^{(2)} A(t)+\ldots+I^{(n)} A(t)+\ldots \\
(I A)(t)=I^{(I)} A(t)=\int_{0}^{t} A(\xi) d \xi \\
\left(I^{(n)} A\right)(t)=\int_{0}^{t} A(\xi)\left(I^{(n-1)} A\right)(\xi) d \xi, n>1
\end{array}\right.
$$

Furthermore, the following theorem is given in [5].
Theorem 3. The equations $\frac{d \psi}{d t}=A(t) \psi$ possess a non-vanishing periodic solution of period $\omega$, if and only if det $(M(\omega))=$ 0. In particular, in order that the equations $\frac{d \psi}{d t}=A(t) \psi$ possess $n$ linearly independent periodic solutions of period $\omega$, the necessary and sufficient condition is that $M(\omega)$ be a zero matrix.

Now, let us apply this theorem to the system (27). If $M(\omega)=0$, there exist the unit vector functions $\mathbf{T}, \mathbf{N}, \mathbf{B}, \mathbf{E}$ of period $\omega$, such that each set of functions $\left\{t_{i}, n_{i}, b_{i}, \varepsilon_{i}\right\},(i=1,2,3,4)$ form a solution of the equation (27) corresponding to the initial conditions $\left(A_{i}, B_{i}, C_{i}, D_{i}\right)$. The curve $(C)$ can be described as follows

$$
\alpha(s)=\int_{0}^{s} \mathbf{T}(s) d s \quad \text { or } \quad \alpha(\varphi)=\int_{0}^{\varphi} \rho(\varphi) \mathbf{T}(\varphi) d(\varphi)
$$

Here, to find $\mathbf{T}$, we can make use of the equation

$$
\left[\begin{array}{c}
t_{i}  \tag{30}\\
n_{i} \\
b_{i} \\
\varepsilon_{i}
\end{array}\right]=\{E+M(\varphi)\}\left[\begin{array}{c}
A_{i} \\
B_{i} \\
C_{i} \\
D_{i}
\end{array}\right],(i=1,2,3,4)
$$

which is established by (29). If we take the initial conditions as $t_{i}(0)=A_{i}, n_{i}(0)=B_{i}, b_{i}(0)=C_{i}, \varepsilon_{i}(0)=D_{i},(i=1,2,3,4)$ such that $\left(A_{1}, A_{2}, A_{3}, A_{4}\right),\left(B_{1}, B_{2}, B_{3}, B_{4}\right),\left(C_{1}, C_{2}, C_{3}, C_{4}\right),\left(D_{1}, D_{2}, D_{3}, D_{4}\right)$ form an orthonormal frame, then from (30) we obtain

$$
\begin{equation*}
t_{i}=\left(1+m_{11}\right) A_{i}+m_{12} B_{i}+m_{13} C_{i}+m_{14} D_{i} ; \quad(i=1,2,3,4) . \tag{31}
\end{equation*}
$$

When the timelike curve $(C)$ is a constant breadth curve, which is also periodic of period $\omega$, it is clear that

$$
\begin{equation*}
\int_{0}^{\omega} \rho t_{i} d \varphi=0 . \tag{32}
\end{equation*}
$$

Hence, form (31) and (32), we have

$$
A_{i} \int_{0}^{\omega} \rho\left(1+m_{11}\right) d \varphi+B_{i} \int_{0}^{\omega} \rho m_{12} d \varphi+C_{i} \int_{0}^{\omega} \rho m_{13} d \varphi+D_{i} \int_{0}^{\omega} \rho m_{14} d \varphi=0 ;(i=1,2,3,4) .
$$

Since the coefficient determinant $\Delta \neq 0$ in this system, we obtain the equalities

$$
\begin{equation*}
\int_{0}^{\omega} \rho\left(1+m_{11}\right) d \varphi=0=\int_{0}^{\omega} \rho m_{12} d \varphi=\int_{0}^{\omega} \rho m_{13} d \varphi=\int_{0}^{\omega} \rho m_{14} d \varphi \tag{33}
\end{equation*}
$$

which are the conditions for a timelike curve to be constant breadth curve in $E_{1}^{4}$. Here, we can take the period $\omega=2 \pi$ because of $0 \leq \varphi \leq 2 \pi$. Thus we establish the following corollary.

Corollary 3. Let $(C)$ be a regular curve in $E_{1}^{4}$ such that $\rho(\varphi)>0, k_{2}(\varphi)$ and $k_{3}(\varphi)$ are continuous periodic functions of period $\omega$. Then $(C)$ is a constant breadth timelike curve and also periodic of period $\omega$, if and only if

$$
\begin{equation*}
M(\omega)=0, \quad \int_{0}^{\omega} \rho\left(1+m_{11}\right) d \varphi=0=\int_{0}^{\omega} \rho m_{12} d \varphi=\int_{0}^{\omega} \rho m_{13} d \varphi=\int_{0}^{\omega} \rho m_{14} d \varphi, \tag{34}
\end{equation*}
$$

holds, where

$$
\left\{\begin{align*}
M(t) & =I A(t)+I^{(2)} A(t)+\ldots+I^{(n)} A(t)+\ldots  \tag{35}\\
A(t) & =\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
1 & 0 & \rho k_{2} & 0 \\
0-\rho k_{2} & 0 & \rho k_{3} \\
0 & 0 & -\rho k_{3} & 0
\end{array}\right]
\end{align*}\right.
$$

and $m_{i j}(t)$ are the entries of the matrix $M(t)$.

By means of (29) and (35), the matrix $M(t)$ can be constructed and each $m_{i j}$ involves infinitely many integrations. Hence, we can write the conditions (34) in the following forms:

$$
\left\{\begin{array}{l}
\int_{0}^{\omega} \rho(\varphi) d \varphi+\int_{0}^{\omega} \int_{0}^{r} \int_{0}^{s} \rho(\varphi) d s d t d \varphi  \tag{36}\\
+\int_{0}^{\omega} \int_{0}^{\phi} \int_{0}^{p} \int_{0}^{r} \int_{0}^{s} \rho(\varphi)[1-\lambda(p) \lambda(s)] d t d s d r d p d \varphi+\ldots=0 \\
\int_{0}^{\omega} \int_{0}^{s} \rho(\varphi) d t d \varphi+\int_{0}^{\omega} \int_{0}^{p} \int_{0}^{r} \int_{0}^{s} \rho(\varphi)[1-\lambda(t) \lambda(s)] d t d s d r d \varphi+\ldots=0 \\
\int_{0}^{\omega} \int_{0}^{r} \int_{0}^{s} \rho(\varphi) \lambda(t) d t d s d \varphi \\
+\int_{0}^{\omega} \int_{0}^{\phi} \int_{0}^{p} \int_{0}^{r} \int_{0}^{s} \rho(\varphi)[\lambda(t)-\lambda(p)\{\lambda(t) \lambda(s)+\mu(t) \mu(s)\}] d t d s d r d p d \varphi+\ldots=0 \\
\int_{0}^{\omega} \int_{0}^{p} \int_{0}^{r} \int_{0}^{s} \rho(\varphi) \lambda(s) \mu(t) d t d s d r d \varphi \\
+\int_{0}^{\omega} \int_{0}^{q} \int_{0}^{\phi} \int_{0}^{p} \int_{0}^{r} \int_{0}^{s} \rho(\varphi) \lambda(p) \mu(t)[1-\lambda(t) \lambda(s)-\mu(t) \mu(s)] d t d s d p d \phi d \varphi+\ldots=0
\end{array}\right.
$$

where $\lambda(\xi)=p(\xi) k_{2}(\xi), \mu(\xi)=p(\xi) k_{3}(\xi)$.

Example 1.Let us consider the special case $\rho=$ const., $k_{2}=$ const. and $k_{3}=$ const. In this case, from (33), we have

$$
\left\{\begin{array}{l}
\omega+\frac{\omega^{3}}{3!}+\left(1-\rho^{2} k_{2}^{2}\right) \frac{\omega^{5}}{5!}+\left(1-\rho^{2} k_{2}^{2}\right)^{2} \frac{\omega^{7}}{7!}+\ldots=0  \tag{37}\\
\frac{\omega^{2}}{2!}+\left(1-\rho^{2} k_{2}^{2}\right) \frac{\omega^{4}}{4!}+\left(1-\rho^{2} k_{2}^{2}\right)^{2} \frac{\omega^{6}}{6!}+\ldots=0 \\
k_{2}\left[\frac{\omega^{3}}{3!}+\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right) \frac{\omega^{5}}{5!}+\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right)^{2} \frac{\omega^{7}}{7!}+\ldots\right]=0 \\
k_{2} k_{3}\left[\frac{\omega^{4}}{4!}+\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right) \frac{\omega^{6}}{6!}+\ldots\right]=0
\end{array}\right.
$$

or

$$
\left\{\begin{array}{l}
\rho^{2} k_{2}^{2}\left(1-\rho^{2} k_{2}^{2}\right)^{\frac{1}{2}} \omega-\sinh \left[\left(1-\rho^{2} k_{2}^{2}\right)^{\frac{1}{2}} \omega\right]=0,  \tag{38}\\
\cosh \left[\left(1-\rho^{2} k_{2}^{2}\right)^{\frac{1}{2}} \omega\right]=1 \text { or }\left(1-\rho^{2} k_{2}^{2}\right)^{\frac{1}{2}} \omega=2 k \pi,(k \in \mathbb{Z}), \\
k_{2}\left[\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right)^{\frac{1}{2}} \omega-\sinh \left[\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right)^{\frac{1}{2}} \omega\right]\right]=0, \\
k_{2} k_{3}\left[\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right) \frac{\omega^{2}}{4}-\sinh ^{2}\left[\left(1-\rho^{2} k_{2}^{2}-\rho^{2} k_{3}^{2}\right)^{\frac{1}{2}} \omega\right]\right]=0,
\end{array}\right.
$$

where $\omega=2 k \pi$. It is seen that all of the equalities (37) or (38) are satisfied simultaneously, if and only if $\rho k_{2}=0$, $\rho k_{3}=0$ that is, $\rho=$ const.$>0$ and $k_{2}, k_{3}=0$. Therefore, only ones with $\rho=$ const. $>0$ and $k_{2}, k_{3}=0$ of the curves with $\rho=$ const.$>0$ and $k_{2}, k_{3}=$ const. are curves of constant breadth, which are Lorentzian circles in $E_{1}^{4}$.

Now let us construct the relation characterizing these circles. Since $\rho k_{2}=1, \rho k_{3}=0$, system (7) becomes

$$
\begin{equation*}
m_{1}^{\prime}=-m_{2}, m_{2}^{\prime}=m_{3}-m_{1}, m_{3}^{\prime}=-m_{2}, m_{4}^{\prime}=0 . \tag{39}
\end{equation*}
$$

The general solution of (39), is

$$
\left\{\begin{array}{l}
m_{1}=\frac{c_{1}}{2} \varphi^{2}+c_{2} \varphi+c_{3}  \tag{40}\\
m_{2}=-c_{1} \varphi-c_{2} \\
m_{3}=\frac{c_{1}}{2} \varphi^{2}+c_{2} \varphi+c_{3}-c_{1} \\
m_{3}=c_{4}
\end{array}\right.
$$

Consequently, replacing (40) into (1), we obtain the equation

$$
\alpha^{*}(\varphi)=\alpha(\varphi)+\left(\frac{c_{1}}{2} \varphi^{2}+c_{2} \varphi+c_{3}\right) \mathbf{T}+\left(-c_{1} \varphi-c_{2}\right) \mathbf{N}+\left(\frac{c_{1}}{2} \varphi^{2}+c_{2} \varphi+c_{3}-c_{1}\right) \mathbf{B}+c_{4} \mathbf{E}
$$

which represents the Lorentzian circles with the diameter

$$
d=\left\|\alpha^{*}-\alpha\right\|=\left(c_{1}^{2}+c_{2}^{2}+c_{4}^{2}-2 c_{1} c_{3}\right)^{\frac{1}{2}}
$$

In this case, a pair of opposite points of the curve is $\left(\alpha^{*}(\varphi), \alpha(\varphi)\right)$ for $\varphi$ in $0 \leq \varphi \leq 2 \pi$.

## 4 Conclusion

In the characterizations and determinations of the special curves and curve pairs are important in the curve theory. A differential equation or a system of differential equations with respect to the curvatures can determinate the special curves or curve pairs. In this paper, the differential equations characterizing the constant breadth timelike curves in are studied $E_{1}^{4}$. Furthermore, a criterion for a timelike space curve to be the curve of constant breadth in $E_{1}^{4}$ is given.

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