

Some new traveling wave solutions of the modified Benjamin-Bona-Mahony equation via the improved (G'/G)-expansion method

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Abstract: The improved (G'/G)-expansion method is a powerful mathematical tool for solving nonlinear evolution equations which arise in mathematical physics, engineering sciences and other technical arena. In this article, we construct some new exact traveling wave solutions for the modified Benjamin-Bona-Mahony equation by applying the improved (G'/G)-expansion method. In the method, the general solution of the second order linear ordinary differential equation with constant coefficients is used for studying nonlinear partial differential equations. The solution procedure of this method is executed by algebraic software, such as, Maple. The obtained solutions including solitary and periodic wave solutions are presented in terms of the hyperbolic function, the trigonometric function and the rational forms. It is noteworthy to reveal that some of our solutions are in good agreement with the published results for special cases which certifies our other solutions. Furthermore, the graphical presentations of some solutions are illustrated in the figures.

Keywords: Traveling wave solutions, the modified Benjamin-Bona-Mahony equation, the improved (G'/G)-expansion method, nonlinear evolution equations.

Keywords: Benjamin-Bona-Mahony equation, traveling wave solution, (G'/G)-expansion method.

1 Introduction

Partial differential equations (PDEs) have been a most important subject of study in all areas of mathematical physics, engineering sciences and other technical arena. At present time, different methods are being established to construct traveling wave solutions of nonlinear PDEs, such as, the tanh method [1], the homogeneous balance method [2], the variational iteration method [3]-[7], the Backlund transformation method [8], the Hirota bilinear transformation method [9], the Jacobi elliptic function expansion method [10,11], the Cole-Hopf transformation method [12], the direct algebraic method [13], the Exp-function method [14]-[19], the Lie symmetry method [20,21], the homotopy analysis method [22] and others [23]-[26].

Recently, Wang et al. [27] presented a method called the (G'/G)-expansion method and obtained exact traveling wave solutions of nonlinear PDEs. In this method, $u(\xi) = \sum_{i=0}^{m} a_i (\frac{G'}{G})^i$, is used, as the traveling wave solutions, where $a_m \neq 0$. After that, many researchers executed this method to establish analytical solutions of the nonlinear PDEs [28]-[41].

More lately, Zhang et al. [42] expanded the basic (G'/G)-expansion method, and called the improved (G'/G)-expansion method for constructing traveling wave solutions of some nonlinear PDEs. In the method, they used,

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 $u(\xi) = \sum_{j=-m}^{m} a_j (\frac{G'}{G})^j$, as the traveling wave solutions, where a_{-m} or a_m may be zero, but both a_{-m} and a_m cannot be zero at a time.

Afterwards, many researchers implemented this powerful and straightforward method to obtain many new traveling wave solutions of various nonlinear PDEs. For instance, Zhao et al. [43] studied the variant Boussinesq equations via this method to construct traveling wave solutions while Hamad et al. [44] concerned about the same method to obtain analytical solutions for the higher dimensional potential YTSF equation. In Ref. [45] Naher and Abdullah studied nonlinear reaction diffusion equation to construct traveling wave solutions via this method whilst they [46] obtained abundant analytical solutions of the (2+1)-dimensional modified Zakharov-Kuznetsov equation by applying the same method. Nofel et al. [47] investigated the fifth order KdV equation by using this method to find exact wave solutions and so on.

Many researchers studied the modified Benjamin-Bona-Mahony equation for obtaining exact traveling wave solutions by using different methods. For example, Yusufoglu and Bekir [48] investigated this equation by applying the tanh and sine-cosine methods to find exact solutions while Yusufoglu [49] implemented the Exp-function method to construct traveling wave solutions of the same equation. In Ref. [50], Taghizadeh and Mirzazadeh used modified extended tanh method to establish traveling wave solutions of this equation. Abbasbandy and Shirzadi [51] executed the first integral method to seek analytical solutions of the same equation. Aslan [52] studied the same equation to find exact solutions by utilizing the basic (G'/G)-expansion method. But, to the best of our knowledge, the modified Benjamin-Bona-Mahony equation is not investigated via the improved (G'/G)-expansion method and the basic-expansion method is that the traveling wave solutions $u(\xi) = \sum_{j=-m}^{m} a_j (\frac{G'}{G})^j$ is used instead of $u(\xi) = \sum_{i=0}^{m} a_i (\frac{G'}{G})^i$.

In this article, we study the partial differential equation, namely, the modified Benjamin-Bona-Mahony equation via the improved (G'/G)-expansion to obtain abundant new exact traveling wave solutions including solitons and periodic solutions.

2 Explanation of the improved -expansion method for the nonlinear PDEs

Suppose the general nonlinear partial differential equation:

$$R(w, w_t, w_x, w_{xx}, w_{xxx}, w_{xt}, ...) = 0,$$
(1)

where w = w(x,t) is an unknown function, R is a polynomial in w = w(x,t) and the subscripts indicate the partial derivatives.

The main steps of the improved (G'/G)-expansion method [42] are:

Step 1. Consider the traveling wave variable:

$$w(x,t) = p(\theta), \quad \theta = Kx + Ht,$$
 (2)

Now using Eq. (2), Eq. (1) is converted into an ordinary differential equation (ODE) for $p(\theta)$:

$$S(p, p', p'', p''', ...) = 0,$$
 (3)

where the superscripts stand for the ordinary derivatives with respect to θ .

Step 2. According to possibility, Eq. (3) can be integrated term by term one or more times, yields constant(s) of

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integration. The integral constant may be zero, for simplicity.

Step 3. Suppose that the traveling wave solution of Eq. (3) can be expressed in the form [42]:

$$p(\boldsymbol{\theta}) = \sum_{j=-n}^{n} r_j (\frac{G'}{G})^j \tag{4}$$

with $G = G(\theta)$ satisfies the second order linear ODE:

$$G'' + \lambda G' + \mu G = 0, \tag{5}$$

where r_j ($j = 0, \pm 1, \pm 2, ..., \pm n$), λ and μ are constants, where either r_{-n} or r_n may be zero, but both r_{-n} and r_n cannot be zero at a time.

Step 4. To determine the value of n, substituting Eq. (4) along with Eq. (5) into Eq. (3) and then taking the homogeneous balance between the highest order derivatives and the highest order nonlinear terms appearing in Eq. (3).

Step 5. Substitute Eq. (4) and Eq. (5) into Eq. (3) with the value of obtained in Step 4. Equating the coefficients of $(\frac{G'}{G})^k$, $(k = 0, \pm 1, \pm 2, ...)$ then setting each coefficient to zero, we obtain a set of algebraic equations for $r_j(j = 0, \pm 1, \pm 2, ..., \pm n)$, H, λ and μ .

Step 6. Solve the system of algebraic equations obtained in step 5 by the help of algebraic software Maple and yields values for $r_j (j = 0, \pm 1, \pm 2, ..., \pm n), H, \lambda$ and μ . Then, substituting obtained values in Eq. (4) along with Eq. (5) with the value of *n*, we obtain traveling wave solutions of the partial differential Eq. (1).

3 Applications of the method

In this section, we study the modified Benjamin-Bona-Mahony equation by applying the improved (G'/G)-expansion method to construct exact traveling wave solutions.

3.1 The modified Benjamin-Bona-Mahony equation

In this work, we consider the modified Benjamin-Bona-Mahony equation followed by

$$u_t + \alpha u_x + \beta u^2 u_x - \gamma u_{xxt} = 0. \tag{6}$$

where λ , γ are free parameters and $\beta \neq 0$.

Now, we use the wave transformation Eq. (2) into Eq. (6), which yields:

$$(H + K\alpha)p' + K\beta p^2 p' - HK^2 \gamma p''' = 0, \tag{7}$$

where the superscripts stand for the derivatives with respect to θ .

Eq. (7) is integrable, therefore, integrating with respect to once yields:

$$(H+K\alpha)p + \frac{K\beta}{3}p^3 - HK^2\gamma p'' + C = 0, \qquad (8)$$

where C is an integral constant which is to be determined later.



Taking the homogeneous balance between p^3 and p'' in Eq. (8), we obtain n = 1. So, the solution of Eq. (8) is:

$$p(\theta) = r_{-1} \left(\frac{G'}{G}\right)^{-1} + r_0 + r_1 \left(\frac{G'}{G}\right)$$
(9)

where r_{-1} , r_0 and r_1 are constants.

Substituting Eq. (9) together with Eq. (5) into the Eq. (8), the left-hand side of Eq. (8) is converted into a polynomial of $\left(\frac{G'}{G}\right)^k$, $(k = 0, \pm 1, \pm 2, ...)$. Collecting all terms with the same power of $\left(\frac{G'}{G}\right)$ according to step 5. Then setting each coefficient of the resulted polynomial to zero, yields a set of algebraic equations for $r_{-1}, r_0, r_1, H, C, \lambda$ and μ (which are not shown, for simplicity).

Solving the system of obtained algebraic equations with the aid of algebraic software Maple, we obtain three different values.

Case 1:

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$$r_{-1} = 0, \qquad r_0 = \pm K\lambda \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^2\gamma(\lambda^2 - 4\mu))}},$$

$$r_1 = \pm 2K \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^2\gamma(\lambda^2 - 4\mu))}}, \qquad H = \frac{-2\alpha K}{2+K^2\gamma(\lambda^2 - 4\mu)}, C = 0$$
(10)

where λ , γ are free parameters and $\beta \neq 0$.

Case 2:

$$r_{-1} = \pm 2K\mu \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^{2}\gamma(\lambda^{2}-4\mu))}}, \quad r_{0} = \pm K\lambda \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^{2}\gamma(\lambda^{2}-4\mu))}},$$

$$r_{1} = 0, \qquad \qquad H = \frac{-2\alpha K}{2+K^{2}\gamma(\lambda^{2}-4\mu)}, C = 0,$$
(11)

where λ, γ are free parameters and $\beta \neq 0$.

Case 3:

$$r_{-1} = \pm 2K\mu \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^2\gamma(\lambda^2+8\mu))}}, \qquad r_0 = \pm K\lambda \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^2\gamma(\lambda^2+8\mu))}},$$
$$r_1 = \pm 2K \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^2\gamma(\lambda^2+8\mu))}}, \qquad H = \frac{-2\alpha K}{2+K^2\gamma(\lambda^2+8\mu)},$$
$$(12)$$
$$= \pm \frac{8K^4\alpha\gamma\mu\lambda}{2+K^2\gamma(\lambda^2+8\mu)} \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^2\gamma(\lambda^2+8\mu))}}$$

where λ, γ are free parameters and $\beta \neq 0$.

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We obtain three different families of traveling wave solutions of Eq. (8) by substituting the general solution Eq. (5) into Eq. (9):

Family 1: Hyperbolic function solutions:

When $\lambda^2 - 4\mu > 0$, we obtain



$$p(\theta) = r_{-1} \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{X \sinh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \cosh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)}{X \cosh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \sinh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)} \right) - \frac{\lambda}{2} \right)^{-1}$$
(13)
+ $r_0 + r_1 \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{X \sinh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \cosh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)}{X \cosh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \sinh(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)} \right) - \frac{\lambda}{2} \right)$

where X and Y are arbitrary constants. If X and Y are taken particular values, various known solutions can be rediscovered.

Family 2: Trigonometric function solutions:

When $\lambda^2 - 4\mu < 0$, we obtain

$$p(\theta) = r_{-1} \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{-X \sin(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \cos(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)}{X \cos(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \sin(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)} \right) - \frac{\lambda}{2} \right)^{-1}$$
(14)
+ $r_0 + r_1 \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{-X \sin(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \cos(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)}{X \cos(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta) + Y \sin(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\theta)} \right) - \frac{\lambda}{2} \right)$

where X and Y are arbitrary constants. If X and Y are taken particular values, various known solutions can be rediscovered.

Family 3: Rational function solutions:

When $\lambda^2 - 4\mu = 0$, we obtain

$$p(\theta) = r_{-1} \left(\frac{Y}{X + Y\theta} - \frac{\lambda}{2} \right)^{-1} + r_0 + r_1 \left(\frac{Y}{X + Y\theta} - \frac{\lambda}{2} \right).$$
(15)

Family 1: Hyperbolic function solutions:

Now, substituting Eqs. (10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), yields the hyperbolic function solution Eq. (13), we obtain following traveling wave solutions respectively (if X = 0 but $Y \neq 0$):

$$p_1(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma(\lambda^2 - 4\mu)}{\beta(2 + K^2\gamma(\lambda^2 - 4\mu))}} \coth \frac{\sqrt{\lambda^2 - 4\mu}}{2} \theta,$$

$$p_2(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{\beta(2 + K^2\gamma(\lambda^2 - 4\mu))}} \left(2\mu \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \coth \frac{\sqrt{\lambda^2 - 4\mu}}{2} \theta - \frac{\lambda}{2} \right)^{-1} + \lambda \right),$$

$$\frac{2\alpha\kappa}{2} = t$$

where $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 - 4\mu)} t$.

$$p_{3}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{\beta(2+\kappa^{2}\gamma(\lambda^{2}-4\mu))}} \left(2\mu \left(\frac{\sqrt{\lambda^{2}-4\mu}}{2} \operatorname{coth} \frac{\sqrt{\lambda^{2}-4\mu}}{2} \theta - \frac{\lambda}{2} \right)^{-1} + \sqrt{\lambda^{2}-4\mu} \operatorname{coth} \frac{\sqrt{\lambda^{2}-4\mu}}{2} \theta \right),$$

where $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma(\lambda^2 + 8\mu)}t$.

Also, substituting Eqs. (10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), yields the



hyperbolic function solution Eq. (13), our traveling wave solutions become respectively (if X = 0 but $Y \neq 0$):

$$p_4(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma(\lambda^2 - 4\mu)}{\beta(2 + K^2\gamma(\lambda^2 - 4\mu))}} \tanh \frac{\sqrt{\lambda^2 - 4\mu}}{2} \theta,$$

$$p_5(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{\beta(2 + K^2\gamma(\lambda^2 - 4\mu))}} \left(2\mu \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \tanh \frac{\sqrt{\lambda^2 - 4\mu}}{2} \theta - \frac{\lambda}{2} \right)^{-1} + \lambda \right)$$

$$p_6(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{\beta(2 + K^2\gamma(\lambda^2 - 4\mu))}} \left(2\mu \left(\frac{\sqrt{\lambda^2 - 4\mu}}{2} \tanh \frac{\sqrt{\lambda^2 - 4\mu}}{2} \theta - \frac{\lambda}{2} \right)^{-1} + \sqrt{\lambda^2 - 4\mu} \tanh \frac{\sqrt{\lambda^2 - 4\mu}}{2} \theta \right),$$

Again, substituting Eqs. (10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), yields the hyperbolic function solution Eq. (13), our obtained exact solutions become respectively. (if $X \neq 0$ but X > Y):

$$p_{7}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma(\lambda^{2}-4\mu)}{\beta(2+K^{2}\gamma(\lambda^{2}-4\mu))}} \tanh\left(\frac{\sqrt{\lambda^{2}-4\mu}}{2}\theta + \theta_{0}\right),$$

$$p_{8}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^{2}\gamma(\lambda^{2}-4\mu))}} \left(2\mu\left(\frac{\sqrt{\lambda^{2}-4\mu}}{2}\tanh\left(\frac{\sqrt{\lambda^{2}-4\mu}}{2}\theta + \theta_{0}\right) - \frac{\lambda}{2}\right)^{-1} + \lambda\right).$$

where $\theta_0 = \tanh^{-1} \frac{Y}{X}$ and $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 - 4\mu)} t$.

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$$p_{9}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{\beta(2+K^{2}\gamma(\lambda^{2}-4\mu))}} \left(2\mu \left(\frac{\sqrt{\lambda^{2}-4\mu}}{2} \tanh\left(\frac{\sqrt{\lambda^{2}-4\mu}}{2} \theta + \theta_{0} \right) - \frac{\lambda}{2} \right)^{-1} + \sqrt{\lambda^{2}-4\mu} \tanh\left(\frac{\sqrt{\lambda^{2}-4\mu}}{2} \theta + \theta_{0} \right) \right),$$

where $\theta_0 = \tanh^{-1} \frac{Y}{X}$ and $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 + 8\mu)} t$.

Family 2: Trigonometric function solutions:

Substituting Eqs. (10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), yields the trigonometric function solution Eq. (14), our exact traveling wave solutions become respectively (if X = 0 but $Y \neq 0$):

$$p_{10}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma(4\mu - \lambda^2)}{\beta(2 + K^2\gamma(4\mu - \lambda^2))}} \cot \frac{\sqrt{4\mu - \lambda^2}}{2} \theta,$$
$$p_{11}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma}{\beta(2 + K^2\gamma(4\mu - \lambda^2))}} \left(2\mu \left(\frac{\sqrt{4\mu - \lambda^2}}{2} \cot \frac{\sqrt{4\mu - \lambda^2}}{2} \theta - \frac{\lambda}{2}\right)^{-1} + \lambda \right),$$

where $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 - 4\mu)} t$.

$$p_{12}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma}{\beta(2+K^2\gamma(4\mu-\lambda^2))}} \begin{pmatrix} 2\mu \left(\frac{\sqrt{4\mu-\lambda^2}}{2}\cot\frac{\sqrt{4\mu-\lambda^2}}{2}\theta - \frac{\lambda}{2}\right)^{-1} \\ +\sqrt{\lambda^2 - 4\mu}\cot\frac{\sqrt{4\mu-\lambda^2}}{2}\theta \end{pmatrix},$$

where $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 + 8\mu)} t$.

Furthermore, substituting Eqs. (10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), yields the



trigonometric function solution Eq. (14), our obtained solutions become respectively (if Y = 0 but $X \neq 0$)

$$p_{13}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma(4\mu-\lambda^2)}{\beta(2+K^2\gamma(4\mu-\lambda^2))}} \tan \frac{\sqrt{4\mu-\lambda^2}}{2}\theta,$$

$$p_{14}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma}{\beta(2+K^2\gamma(4\mu-\lambda^2))}} \left(2\mu \left(\frac{\sqrt{4\mu-\lambda^2}}{2}\tan \frac{\sqrt{4\mu-\lambda^2}}{2}\theta - \frac{\lambda}{2}\right)^{-1} + \lambda \right)$$

$$p_{15}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma}{\beta(2+K^2\gamma(4\mu-\lambda^2))}} \left(2\mu \left(\frac{\sqrt{4\mu-\lambda^2}}{2}\tan \frac{\sqrt{4\mu-\lambda^2}}{2}\theta - \frac{\lambda}{2}\right)^{-1} + \sqrt{\lambda^2 - 4\mu}\tan \frac{\sqrt{4\mu-\lambda^2}}{2}\theta} \right),$$

Moreover, substituting Eqs.(10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), yields the trigonometric function solution Eq. (14), we obtain following solutions respectively (if $X \neq 0$ but X > Y):

$$p_{16}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma(4\mu-\lambda^2)}{\beta(2+K^2\gamma(4\mu-\lambda^2))}} \tan\left(\frac{\sqrt{4\mu-\lambda^2}}{2}\theta - \theta_0\right),$$

$$p_{17}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma}{\beta(2+K^2\gamma(4\mu-\lambda^2))}} \left(2\mu\left(\frac{\sqrt{4\mu-\lambda^2\mu}}{2}\tan\left(\frac{\sqrt{4\mu-\lambda^2\mu}}{2}\theta - \theta_0\right) - \frac{\lambda}{2}\right)^{-1} + \lambda\right),$$

$$p_{18}(x,t) = \pm K \sqrt{\frac{3\alpha\gamma}{\beta(2+K^2\gamma(4\mu-\lambda^2\mu))}} \left(2\mu\left(\frac{\sqrt{4\mu-\lambda^2}}{2}\tanh\left(\frac{\sqrt{4\mu-\lambda^2}}{2}\theta - \theta_0\right) - \frac{\lambda}{2}\right)^{-1} + \sqrt{4\mu-\lambda^2}\tanh\left(\frac{\sqrt{4\mu-\lambda^2}}{2}\theta - \theta_0\right)\right),$$

Family 3: Rational function solutions:

Substituting Eqs. (10), (11) and (12) together with the general solution Eq. (5) into the Eq. (9), we obtain rational function solution Eq. (14), and our obtained exact solutions become respectively (if $\lambda^2 - 4\mu = 0$):

$$p_{19}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{2\beta}} + \frac{2Y}{X+Y\theta},$$
$$p_{20}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{2\beta}} \left(2\mu \left(\frac{Y}{X+Y\theta} - \frac{\lambda}{2} \right)^{-1} + \lambda \right),$$

where $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 - 4\mu)} t$.

$$p_{21}(x,t) = \pm K \sqrt{\frac{-3\alpha\gamma}{2\beta}} \left(2\mu \left(\frac{Y}{X+Y\theta} - \frac{\lambda}{2} \right)^{-1} + \frac{2Y}{X+Y\theta} \right),$$

where $\theta = Kx - \frac{2\alpha K}{2 + K^2 \gamma (\lambda^2 + 8\mu)} t$.

4 Results and discussion

It is important to state that some of our solutions are in good agreement with the published results and some are new in this article. In addition, the graphical presentations of some obtained solutions are described in the following subsection in figure 1 to figure 12.

Aslan [52] solutions	Our solutions
i. If $C_1 \neq 0, C_2=0, \lambda=3, \mu=2,$	i. If $\lambda = 3, \mu = 2, \alpha = \beta = \gamma = 1$,
$\alpha = \beta = \gamma = 1$ and $k = 1$, solution (33)	$K = 1$ and $p_1 = u_{1,2}$
(from: subsection 3.3) becomes:	solution $p_1(\theta)$ becomes:
$u_{1,2}(x,t) = \pm i \coth \frac{1}{2}(x - \frac{2}{3}t).$	$u_{1,2}(x,t) = \pm i \coth \frac{1}{2}(x - \frac{2}{3}t).$
ii. If $C_1=0, C_2 \neq 0, \lambda=3, \mu=2,$	ii. If $\lambda = 3, \mu = 2, \alpha = \beta = \gamma = 1$,
$\alpha = \beta = \gamma = 1$ and $k = 1$, solution (33)	$K = 1$ and $p_4 = u_{1,2}$
(from: subsection 3.3) becomes:	solution $p_1(\theta)$ becomes:
$u_{1,2}(x,t) = \pm i \tanh \frac{1}{2}(x - \frac{2}{3}t).$	$u_{1,2}(x,t) = \pm i \tanh \frac{1}{2}(x - \frac{2}{3}t).$
iii. If $C_1=0, C_2 \neq 0, \lambda=4, \mu=5,$	iii. If $\lambda = 4, \mu = 5, \alpha = \beta = \gamma = 1,$
$\alpha = \beta = \gamma = 1$ and $k = 1$, solution (34)	$K = 1$ and $p_{10} = u_{3,4}$
(from: subsection 3.3) becomes:	solution $p_{10}(\theta)$ becomes:
$u_{3,4}(x,t) = \pm \sqrt{2} \cot i(x - \frac{1}{3}t).$	$u_{3,4}(x,t) = \pm \sqrt{2} \cot i(x - \frac{1}{3}t).$
iv. If $C_1 \neq 0, C_2=0, \lambda=4, \mu=5$,	iv. If $\lambda = 4, \mu = 5, \alpha = \beta = \gamma = 1,$
$\alpha = \beta = \gamma = 1$ and $k = 1$, solution (34)	$K = 1$ and $p_{13} = u_{3,4}$
(from: subsection 3.3) becomes:	solution $p_{13}(\theta)$ becomes:
$u_{3,4}(x,t) = \pm \sqrt{2} \tan i(x - \frac{1}{3}t).$	$u_{3,4}(x,t) = \pm \sqrt{2} \tan i(x - \frac{1}{3}t).$
v. If $C_3=1, C_2=0, \lambda^2-4\mu=0,$	v. If $X=0, Y=1, \lambda^2-4\mu=0, \alpha=\gamma=1,$
$\alpha = \beta = \gamma = 1$ and $k = 1$, solution (35)	$\beta = 6, K = 1 \text{ and } p_{19} = u_{5,6}$
(from: subsection 3.3) becomes:	solution $p_{19}(\theta)$ becomes:
$u_{5,6}(x,t) = \pm \frac{1}{x-t}$).	$u_{5,6}(x,t) = \pm \frac{1}{x-t}$).
vi. If $\lambda = 3, \mu = 2, \alpha = 3, \beta = -3, \gamma = 1$,	vi. If $\lambda = 3, \mu = 2, \alpha = 3, \beta = -3, \gamma = 1,$
$\xi_0 = 0$ and $k = 1$, solution (36)	$K = 1, \theta_0 = 0$ and $p_7 = u_{7,8}$
(from: subsection 3.3) becomes:	solution $p_7(\theta)$ becomes:
$u_{7,8}(x,t) = \pm \tanh \frac{1}{2}(x-2t).$	$u_{7,8}(x,t) = \pm \tanh \frac{1}{2}(x-2t).$

4.1 Table. Comparison between Aslan [52] solutions and our solutions:

Beside this table, we obtain new traveling wave solutions p_2 , p_3 , p_5 , p_6 , p_8 , p_9 , p_{11} , p_{12} , p_{14} to p_{18} , p_{20} and p_{21} which have not been found in the previous literature.

4.2 Graphical representations of the solutions

Some of our solutions are shown in the figures by using the commercial software Maple:



Fig. 1: Periodic solution for
$$\label{eq:lambda} \begin{split} \lambda = 4, \mu = 6, \alpha = 2, \beta = 3, \gamma = 4, K = 7 \end{split}$$

Fig. 2: Soliton solution for

 $\hat{\lambda} = 4, \mu = 5, \alpha = 2, \beta = 2, \gamma = 1, K = 3$







 $\lambda = 5, \mu = 7, \alpha = 3, \beta = 3, \gamma = 2, K = 5$





 $\hat{\lambda}=6, \mu=10, \alpha=1, \beta=3, \gamma=3, K=4$



Fig. 5: Solitons solution for

 $\hat{\lambda} = 4, \mu = 4, \alpha = 1, \beta = 6, \gamma = 1, K = 5, X = 1, Y = 1$



Fig. 7: Solitons solution for

 $\lambda=4, \mu=4, \alpha=2, \beta=6, \gamma=3, K=8, X=1, Y=1$



Fig. 6: Solitons solution for

 $\hat{\lambda} = 6, \mu = 9, \alpha = 1, \beta = 6, \gamma = 1, K = 1, X = 0, Y = 1$





 $\hat{\lambda} = 4, \mu = 4, \alpha = 2, \beta = 3, \gamma = 1, K = 1, X = 1, Y = 1$



0.4 0.2 0

0.4



Fig. 9: Solitons solution for

 $\hat{\lambda} = 4, \mu = 3, \alpha = 3, \beta = 2, \gamma = 5, K = 1$



 $\hat{\lambda} = 6, \mu = 6, \alpha = 0.5, \beta = 5, \gamma = 7, K = 15$





 Fig. 11: Solitons solution for
 Fig. 12: Solitons solution for

 $\lambda = 3, \mu = 4, \alpha = 1, \beta = 1, \gamma = 1, K = 1$ $\lambda = 8, \mu = 16, \alpha = 2, \beta = 0.5, \gamma = 0.5, K = 0.5, X = 0, Y = 1$

5 Conlusion

In this article, we obtain abundant exact traveling wave solutions of the modified Benjamin-Bona-Mahony equation via the improved (G'/G)-expansion method. The obtained solutions including solitons and periodic solutions are presented in terms of the hyperbolic, the trigonometric and the rational forms. Further, it is important mentioning that some of our solutions are in good agreement with the published results and some are new which validates our other solutions. Moreover, the solutions show that the used method is more effective and powerful mathematical tool and gives many new solutions. We expect that nonlinear PDEs which are arising in mathematical physics, engineering sciences and other technical arena will be investigated by applying this straightforward method.

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